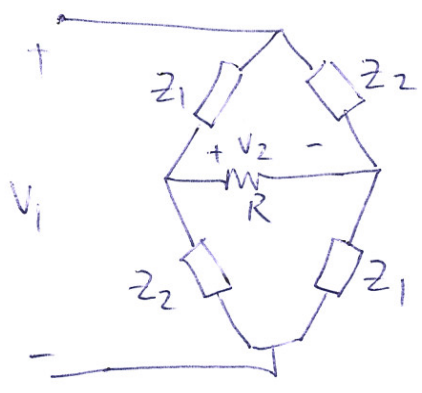
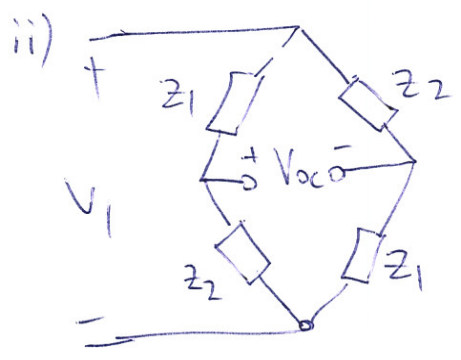


2nd Order All-Pass Circuit:

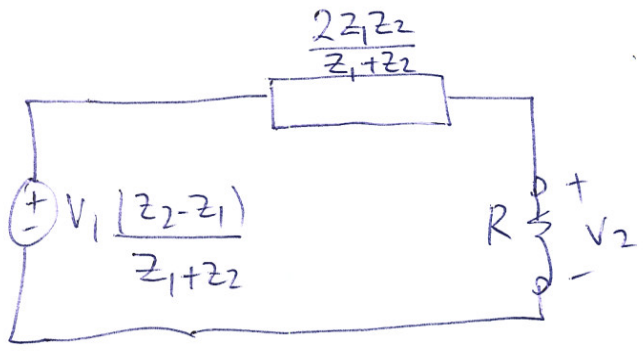


i) Z_{Th} seen by $R \Rightarrow (z_1 || z_2) 2$.



$$V_{oc} = V_1 \left(\frac{z_2 - z_1}{z_1 + z_2} \right)$$

Then,



$$\frac{V_2(s)}{V_1(s)} = ?$$

$$\frac{V_2(s)}{V_1(s)} = \frac{(z_2^{(s)} - z_1^{(s)})}{z_1^{(s)} + z_2^{(s)}} \cdot \frac{R}{R + \frac{2z_1z_2}{z_1+z_2}} = (z_2 - z_1) \cdot \frac{R}{R(z_1+z_2) + 2z_1z_2}$$

$$= \frac{(z_2 - z_1)}{(z_1 + z_2) + \frac{2z_1z_2}{R}} \rightarrow$$

\rightarrow det $z_1 z_2 = R^2$ \rightarrow μ

$$= \frac{(z_2 - z_1)}{(z_1 + z_2) + 2R}$$

$$= \frac{z_2(z_2 - z_1)}{z_2 \{ (z_1 + z_2) + 2R \}}$$

$$H(s) = \frac{z_2 - R}{z_2 + R}$$

$$= \frac{(z_2 - R)(z_2 + R)}{(z_2 + R)^2} = \frac{z_2^2 - R^2}{z_2^2 + 2Rz_2 + R^2}$$

Then

$$H(j\omega) = \frac{Z_2(j\omega) - R}{Z_2(j\omega) + R}$$

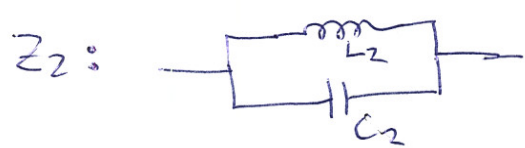
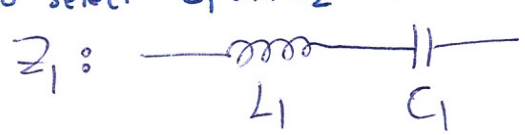
Then ① $Z_2(j\omega)$: purely real $\forall \omega \rightarrow Z_2 = R_2 \rightarrow H(j\omega) = \frac{R_2 - R}{R_2 + R}$

$\rightarrow |H(j\omega)|$ is not function of $\omega \rightarrow$ but system does not contain any dynamic elements. Therefore it is just a voltage divider, not a filter.

② $Z_2(j\omega)$: purely imaginary $\forall \omega \rightarrow Z_2 = jX_2 \rightarrow H(j\omega) = \frac{jX_2 - R}{jX_2 + R}$

$\rightarrow |H(j\omega)| = 1 \rightarrow$ we have dynamic system with an all-pass structure. (Q: If $Z_2 = jX_2$, can $Z_1 Z_2 = R^2$ be satisfied? (No!))

Example: How to select Z_1 and Z_2 such that $Z_1 Z_2 = R^2$ is satisfied?



$$\left\{ \begin{aligned} Z_1(s) &= sL_1 + 1/sC_1 \\ Z_2(s) &= \left(\frac{1}{sL_2} + sC_2 \right)^{-1} \end{aligned} \right.$$

$$Z_1 \cdot Z_2 = R^2 \rightarrow Z_1 = \frac{R^2}{Z_2} \rightarrow sL_1 + \frac{1}{sC_1} = R^2 \left(\frac{1}{sL_2} + sC_2 \right)$$

det

$$\begin{aligned} L_1 &= R^2 C_2 \\ C_1 &= L_2 / R^2 \end{aligned}$$

Then $H(s) = \frac{Z_2 - R}{Z_2 + R} = \frac{1 - RY_2}{1 + RY_2}$ $\left(Z_2^{-1} = Y_2 = \frac{1}{sL_2} + sC_2 \right)$

$$= \frac{1 - \frac{R}{sL_2} - sRC_2}{1 + \frac{R}{sL_2} + sRC_2}$$

$$= - \frac{s^2 - \frac{1}{RC_2}s + \frac{1}{L_2C_2}}{s^2 + \frac{1}{RC_2}s + \frac{1}{L_2C_2}}$$

$\frac{1}{RC_2} \rightarrow 2\alpha\omega_0$ $\frac{1}{L_2C_2} \rightarrow \omega_0^2$

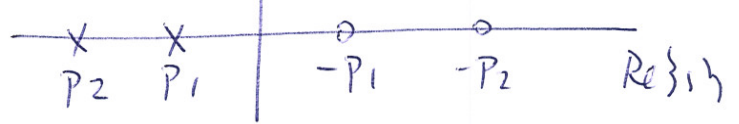
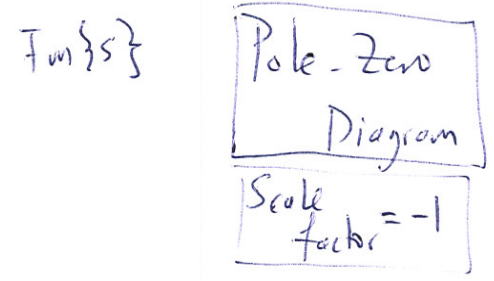
$$H(s) = - \frac{s^2 - 2\alpha\omega_0s + \omega_0^2}{s^2 + 2\alpha\omega_0s + \omega_0^2}$$

Note: If s_1, s_2 are roots of $s^2 + 2\alpha\omega_0s + \omega_0^2 = 0$.

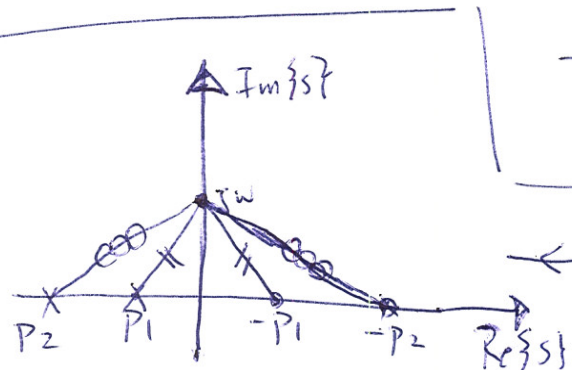
then $-s_1, -s_2$ are the roots of $s^2 - 2\alpha\omega_0s + \omega_0^2 = 0$.
 (Why? Hint: $2\alpha\omega_0 = \text{sum of roots}$; $\omega_0^2 = \text{product of roots}$)

then.

$$H(s) = - \frac{(s - p_1)(s - p_2)}{(s + p_1)(s + p_2)}$$

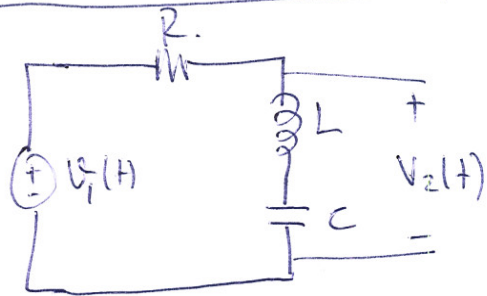


Note:



← This picture shows that system is all-pass!

Series RLC Band-Stop Filter



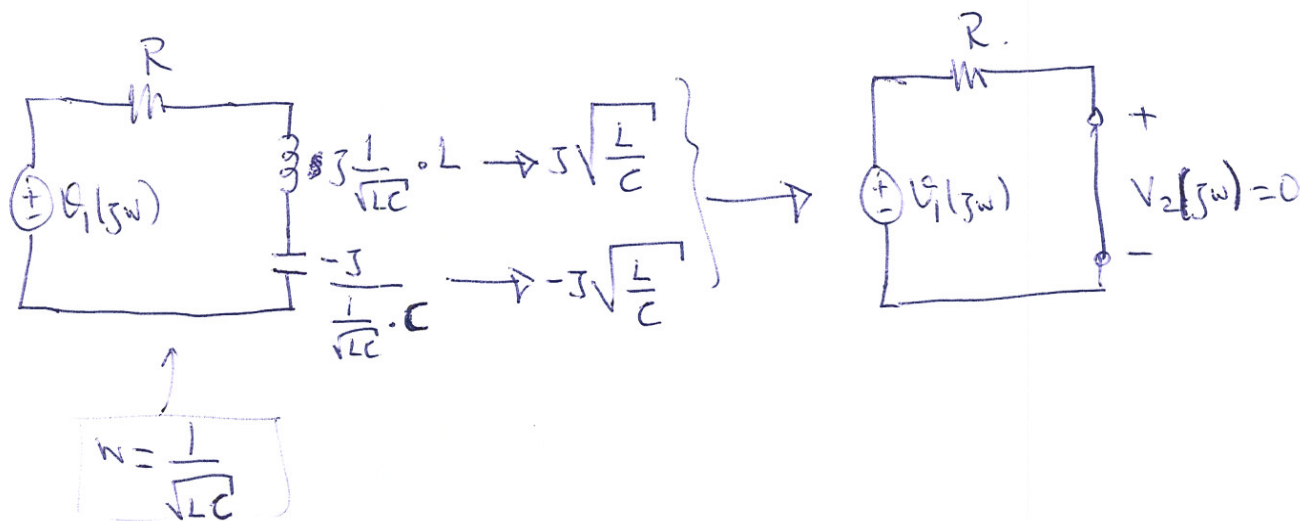
$$H(s) = \frac{V_2(s)}{V_1(s)} = \frac{sL + 1/sC}{sL + 1/sC + R}$$

$$= \frac{s^2 + 1/LC}{s^2 + \frac{R}{L}s + 1/LC}$$

$$H(j\omega) = \frac{(1/LC - \omega^2)}{(1/LC - \omega^2) + j\omega \frac{R}{L}}$$

→ Note: $|H(j\omega)| = 0$ for $\omega = \frac{1}{\sqrt{LC}}$

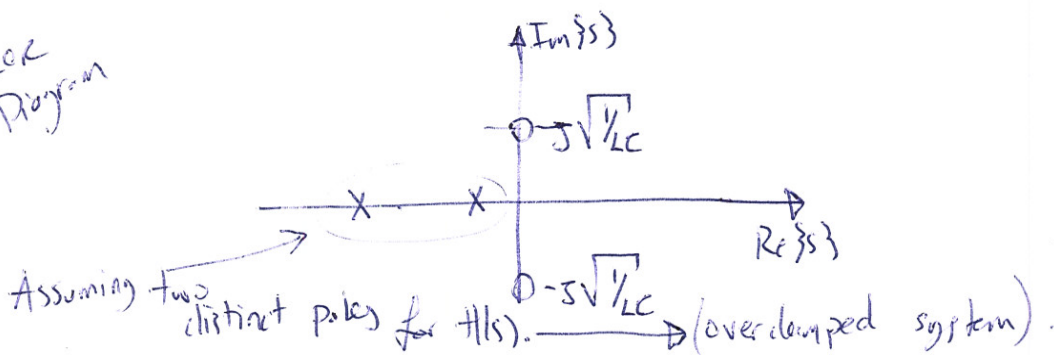
Let's redraw the circuit at $\omega = 1/\sqrt{LC}$



Note that $V_1(j\omega)$ source sees purely resistive circuit at $\omega = 1/\sqrt{LC}$

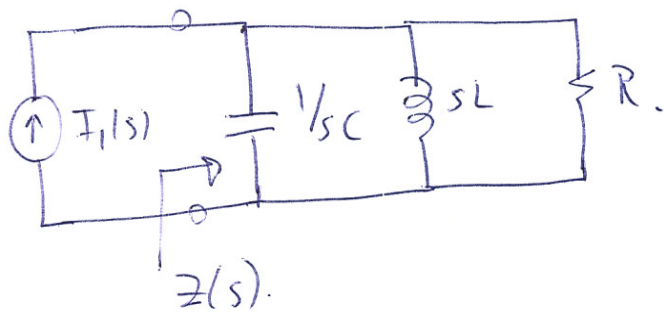
This phenomena is called resonance. (more on this later).

Pole-Zero Diagram



Parallel RLC Resonance Circuit:

3A



$$Z(s) = \frac{1}{\frac{1}{R} + \frac{1}{sL} + sC} = \frac{s/C}{s^2 + \frac{1}{RC}s + \frac{1}{LC}}$$

Definition: The frequency for which $Z(j\omega)$ is purely real is called the resonance frequency.

$$Z(j\omega) = \frac{j\omega/C}{- \omega^2 + \frac{1}{LC} + j\omega \frac{1}{RC}} = \frac{\omega/C}{j(\omega^2 - 1/LC) + \omega/RC}$$

is zero when $\omega = \frac{1}{\sqrt{LC}}$

Resonance freq.

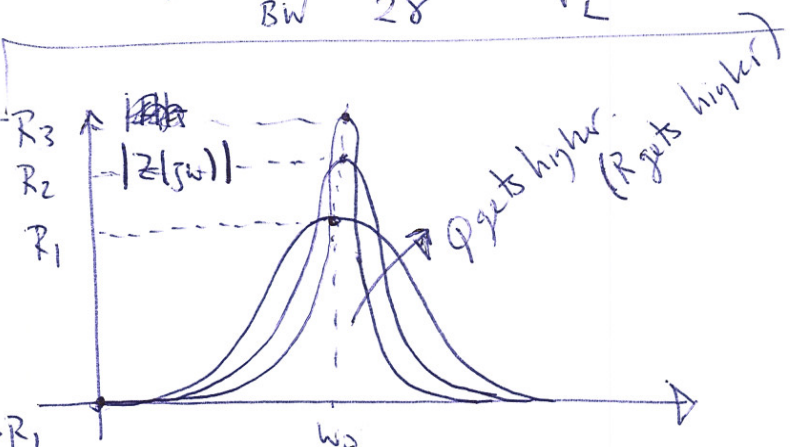
(We have studied this system under the title of 2nd order Band-Pass circuits)

$$s^2 + \frac{1}{RC}s + \frac{1}{LC} \rightarrow \text{BW} = \omega_2 - \omega_1 = 2\gamma\omega_0 = \frac{1}{RC}$$

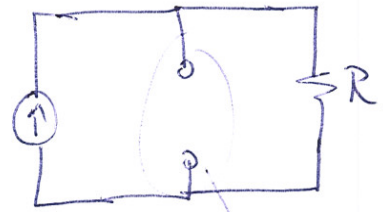
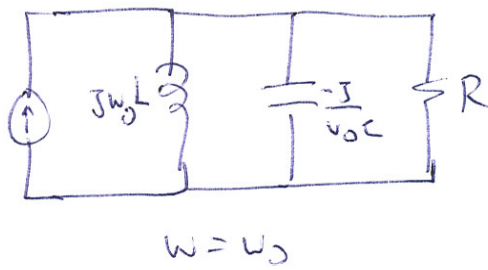
$\omega_0 = \frac{1}{\sqrt{LC}}$ $2\gamma = \frac{1}{R}\sqrt{\frac{L}{C}}$

$$Q = \frac{\omega_0}{\text{BW}} = \frac{1}{2\gamma} = R\sqrt{\frac{C}{L}}$$

Picture shows the impedance seen by the source for 3 parallel RLC circuits with $R_3 > R_2 > R_1$ and the same L and C's.



The maximum impedance is seen at resonance freq. (for parallel RLC) and it is equal to R . Note that (3B)



combination of L and C at $\omega = \omega_0$.

Another Interpretation for Q : Quality factor.

For parallel RLC circuit, the average energy stored in capacitor is $E_C = \frac{1}{2} C V_{eff}^2$; similarly the average magnetic energy stored in inductor is $E_L = \frac{1}{2} L I_{eff}^2$ →

Let's calculate the total energy stored in L and C .

$$E_T = E_C + E_L = \frac{1}{2} C V_{eff}^2 + \frac{1}{2} L \left(\frac{V_{eff}}{\omega L} \right)^2$$

$$= \frac{1}{2} \left(C + \frac{1}{\omega^2 L} \right) V_{eff}^2$$

$$= \frac{1}{2} C \left(1 + \frac{1}{\omega^2 LC} \right) V_{eff}^2$$

$$= \frac{1}{2} C \left(1 + \frac{\omega_0^2}{\omega^2} \right) V_{eff}^2$$

$$= \frac{1}{2} C \frac{\omega_0}{\omega} \left(\frac{\omega}{\omega_0} + \frac{\omega_0}{\omega} \right) V_{eff}^2$$

We will now show that at $\omega = \omega_0$

$$Q = 2\pi \cdot \frac{E_T}{T P_R} \quad \left(\begin{array}{l} Q: \\ \text{unitless} \end{array} \right) \quad \text{where} \quad T = \frac{2\pi}{\omega_0}$$

(4A)

where P_R is the average ~~power~~ ^{power} consumed by R and $T P_R$ is the energy dissipated over R in a period.

Now

$$\begin{aligned} 2\pi \frac{E_T}{T P_R} &= \frac{\omega E_T}{P_R} = \frac{\omega \frac{1}{2} C \frac{\omega_0}{\omega} \left(\frac{\omega}{\omega_0} + \frac{\omega_0}{\omega} \right) V_{\text{eff}}^2}{\frac{V_{\text{eff}}^2}{R}} \\ &= \frac{\cancel{\omega} \frac{1}{2} C \omega_0 \left(\frac{\omega}{\omega_0} + \frac{\omega_0}{\omega} \right)}{R} \\ &= \frac{RC \omega_0}{2} \left(\frac{\omega}{\omega_0} + \frac{\omega_0}{\omega} \right) \\ &= \frac{R \sqrt{\frac{C}{L}}}{2} \left(\frac{\omega}{\omega_0} + \frac{\omega_0}{\omega} \right) \\ &= \frac{Q}{2} \left(\frac{\omega}{\omega_0} + \frac{\omega_0}{\omega} \right) \end{aligned}$$

for $\omega = \omega_0 \rightarrow 2\pi \frac{E_T}{T P_R} = Q, \quad T = \frac{2\pi}{\omega_0}$

Note: Resonant freq. is the freq. where maximum of $|Z(j\omega)|$ for parallel RLC circuit; BUT this is not true in general. If ω_m is the freq. of maxima of $|Z(j\omega)|$ $\omega_m \neq \omega_0$ in general.

(5A)

$$Z(s) = \frac{j\omega/C + r_1/L}{s = j\omega \left(\frac{1}{LC} - \omega^2 \right) + j\omega \frac{r_1}{L}}$$

$$= \frac{1}{C} \frac{(j\omega + \frac{r_1}{L})}{(\omega_a^2 - \omega^2) + j\omega \frac{\omega_a}{Q}}$$

$$= \frac{1}{C} \frac{(j\omega + \frac{\omega_a}{Q})}{(\omega_a^2 - \omega^2) + j\omega \frac{\omega_a}{Q}}$$

for $Z(j\omega)$ to be purely real, angle of numerator and denominator of $Z(j\omega)$ should be same \rightarrow So

$$\frac{\omega}{\omega_a/Q} = \frac{\omega_a/Q}{\omega_a^2 - \omega^2} \rightarrow \left(\frac{\omega_a}{Q} \right)^2 = \omega_a^2 - \omega^2$$

$$\omega^2 = \omega_a^2 \left(1 - \frac{1}{Q^2} \right)$$

$$\omega = \omega_a \sqrt{1 - \frac{1}{Q^2}}$$

then $\omega_0 = \omega_a \sqrt{1 - \frac{1}{Q^2}}$ is the resonant freq.

$$Z(j\omega_0) = \frac{1}{C} \frac{(j\omega_0 + \frac{\omega_a}{Q})}{\frac{\omega_a}{Q} (j\omega_0 + \frac{\omega_a}{Q})} = \frac{Q \cdot 1}{C \omega_a} = \frac{Q \cdot \omega_a^2 L C}{\omega_a}$$

$$= Q \cdot \frac{\omega_a r_1}{r_1/L} = Q^2 \cdot r_1$$

Ex: a) $r_1 = 50 \Omega$, $L = 10 \text{ mH}$, $C = 1 \mu\text{F}$

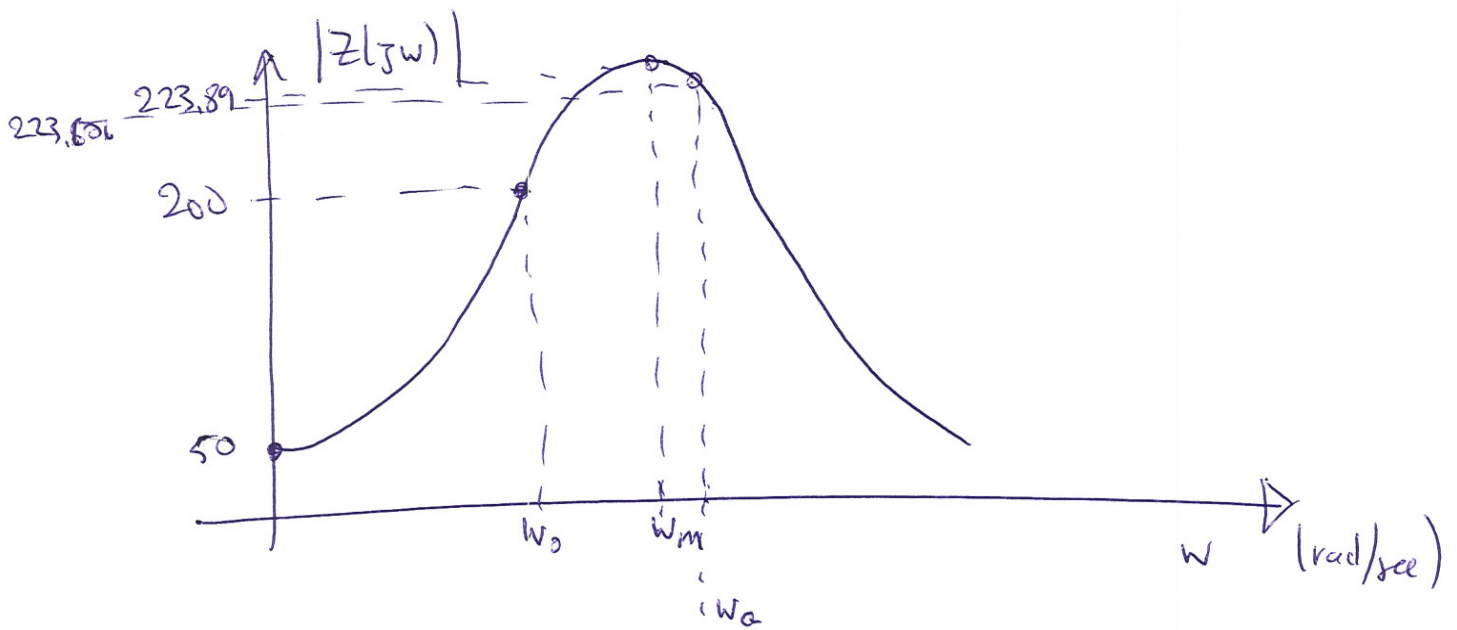
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then

$$\left\{ \begin{aligned} \omega_a &= \frac{1}{\sqrt{LC}} = 10^4 = 10 \text{ k rad/sec.} \\ \phi &= \frac{\omega_a}{r_1/L} = 2 \text{ (unitless).} \\ \omega_0 &= \omega_a \sqrt{1 - \frac{1}{\phi^2}} = 10^4 \sqrt{1 - \frac{1}{4}} = 8.66 \text{ k rad/sec.} \end{aligned} \right.$$

$$Z(s_0) = 50 \Omega, \quad Z(j\omega_0) = \phi^2 \cdot r_1 = 200 \Omega; \quad Z(j\omega_a) = 223.606 \angle -26^\circ$$

$$Z(j\omega_m) = 223.89 \angle -23^\circ$$

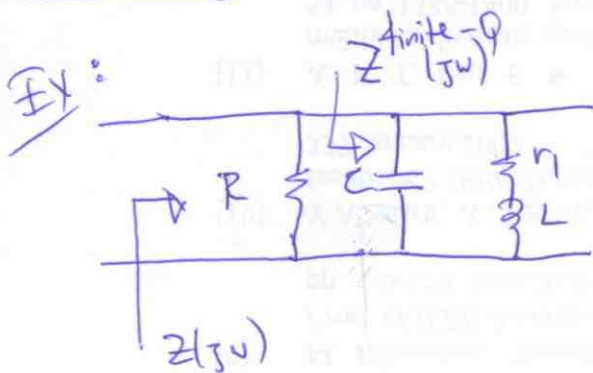


$$\omega_a = \sqrt{1/LC}$$

ω_m : \rightarrow should be found by taking derivative of $|Z(j\omega)|$ wrt. ω .

$\omega_0 = \omega_a \sqrt{1 - 1/\phi^2}$ \leftarrow resonant freq. of finite- ϕ circuit.

We observe that for $Q=2$; ω_a is sufficiently close to ω_m . As $Q \rightarrow \infty$, ω_a and ω_m approaches ω_0 . But note that $Q=2$ is close enough for many purposes in this example. (11) (6A)



Find $Z(j\omega)$, and assume that $\omega \approx$ resonant freq of this system.

