# ODTÜ METU Insights from Swift J0243.6+6124 during its 2017-2018 outburst EBERHARD KARLS UNIVERSITÄT TÜBINGEN



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 $\checkmark$  Uncovered a new transition at  $L_{coul}$ , consistent with B ~  $5 \times 10^{12}$  G and d = 5.2 kpc.  $\checkmark$  Found elevated timing noise strengths above super-Eddington levels, may originate from emerging quadrupole fields.

### Source

- Discovered in 2017. Outbursts in 2017-2018 & 2023.
- Be X-ray binary with  $P_{spin} \sim 9.8$  sec,  $P_{orbit} \sim 27.7$  d <sup>[a][b]</sup>
- Studied in detail: 25+ articles so far!
- 2017-18 outburst: X-ray luminosity varying by 5 orders of magnitude!
- First detected ultraluminous X-ray pulsar (ULX) in the Milky Way! L<sub>peak</sub> ~ 1×10<sup>39</sup> erg/s at 5.2 kpc <sup>[c]</sup>
- Highest-energy CRSF ~ 120-146 keV  $\rightarrow$  B ~1.6×10<sup>13</sup> G, too high! Associated with multipole fields? <sup>[c]</sup>

### **Pulse Timing**

• Strong spin-up at the outburst beginning  $\rightarrow$  phase-coherent timing technique unfavorable. So, we used this approach:

## **Timing Noise**

1 When torque fluctuations are...

• We used the Gaia EDR3

distance:  $5.2 \pm 0.3$  kpc (revised

from 6.8 kpc by Gaia EDR2)

...uncorrelated, wind accretion  $\rightarrow$  white noise (flat,  $\Gamma = 0$ )

### Data

- 480 ks NICER/XTI observations: MJD 58030-58530
- Public Fermi/GBM pulse frequency & Swift/BAT daily light curve histories



Figure 1: Spin frequencies, timing solutions and corresponding *NICER* phase residuals

Previously, 2 transitional luminosity levels reported: L<sub>1</sub> & L<sub>2</sub> <sup>[c]</sup>



- ...correlated, disk accretion  $\rightarrow$  red noise ( $\Gamma$  = -2)
- Generated two power density spectra (PDS) of frequency derivatives via the rms-value technique <sup>[f]</sup> (Figure 3)
- Along with the standard method, we applied a torqueluminosity model to spin frequencies and then used the residual frequencies  $\rightarrow$  minimizes the disk accretion contribution to noise levels
- Standard method: Red noise component has  $\Gamma = -3.36 \pm 0.64$ , steeper than similar sources
- Modified method:  $\Gamma = -0.91 \pm 0.38$ , luminosity-dependent model removed most (but not all) the red noise component



Figure 3: PDS of spin frequency derivatives with broken power law fits

Figure 4: Luminosity dependence of timing noise strengths

- We also checked the luminosity dependence of timing noise strengths using the standard method (Figure 4)
- At the highest luminosities, torque interactions become less efficient & noisier  $\rightarrow$  supports the previous deductions of interactions with quadrupole components of magnetic field [c]

#### • L ~ $7 \times 10^{36}$ erg/s $\rightarrow$ Consistent with L<sub>coul</sub> of Becker's model <sup>[d]</sup>: pencil beam to mixed pencil & fan beam

 $L \rightarrow B \sim 4.7 \times 10^{12} \text{ G}$  all consistent  $L_1 \rightarrow B \sim 5.3 \times 10^{12} \text{ G}$  at d = 5.2 kpc!

**1** Below L<sub>coul</sub>, Coulomb interactions cannot stop the accretion flow: gas shock only.

more  $\rightarrow$  more luminosity **1** Torque-Luminosity Model: accretion  $\rightarrow$  more torque on NS  $\Rightarrow$  more spin-up/down!

• Ghosh-Lamb model <sup>[g]</sup>:  $\dot{v} \propto L^{6/7}$  (We used  $\dot{v}_{model} = \beta L^{\alpha} + \dot{v}_0$ )

## References

[a] Kennea et al. 2017, <u>ATel</u> #10809 [b] gammaray.nsstc.nasa.gov/qbm/science/ pulsars/lightcurves/swiftj0243.html [c] Doroshenko et al. 2020, MNRAS, 491, 1857

[d] Becker et al. 2012, <u>A&A</u>, 544, A123 [e] Kong et al. 2022, ApJ, 933, L3 [f] Deeter 1984, <u>Apl</u>, 281, 482 [g] Ghosh & Lamb 1979, <u>ApJ</u>, 234, 296

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