

STABILITY OF PERIODIC SOLUTIONS OF DIFFERENTIAL EQUATIONS
WITH IMPULSE ACTION ON SURFACES

M. U. Akhmetov and N. A. Perestyuk

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We will consider a system of differential equations with the impulse action

$$dx/dt = f(t, x), \quad t \neq t_i(x), \quad \Delta x|_{t=t_i(x)} = I_i(x), \quad (1)$$

in which $x \in \mathbb{R}^n$, and the functions $f(t, x)$, $I_i(x)$, $t_i(x)$ are defined and continuous in t and x , $t \in \mathbb{R}$, $x \in \mathbb{R}^n$, $i \in \mathbb{Z}$. The function $f(t, x)$ is T -periodic in t , and there exists a number $p \in \mathbb{N}$, for which, for all $x \in \mathbb{R}^n$, $i \in \mathbb{Z}$ we have $I_{i+p}(x) = I_i(x)$, $t_{i+p}(x) = t_i(x) + T$.

The surfaces of the discontinuity satisfy a uniqueness condition which consists in the fact that a solution of Eqs. (1) intersects each of these surfaces no more than once, and also $t_i(x) > t_{i-1}(x)$, $x \in \mathbb{R}^n$, $i \in \mathbb{Z}$.

The question of the asymptotic stability of the family of periodic solutions of (1) depending on $\gamma \in \mathbb{R}^m$, is analyzed here. Beforehand, a theorem was proved on the differentiability with respect to the parameter of the solutions of the system

$$dx/dt = f(t, x, \mu), \quad t \neq t_i(x, \mu), \\ \Delta x|_{t=t_i(x, \mu)} = I_i(x, \mu), \quad (2)$$

where the functions $f(t, x, \mu)$, $I_i(x, \mu)$, and $t_i(x, \mu)$ are defined and continuous for all $t \in \mathbb{R}$, $x \in \mathbb{R}^n$, $\mu \in \mathbb{R}^m$, $i \in \mathbb{Z}$ with respect to t , x , and μ ; the surface $t = t_i(x, \mu)$ satisfy a uniqueness condition and $t_i(x, \mu) > t_{i-1}(x, \mu)$ for $x \in \mathbb{R}^n$, $\mu \in \mathbb{R}^m$, $i \in \mathbb{Z}$.

Differential equations with impulse action arose as mathematical models in the solution of practical problems of importance and now the theory of impulse systems finds wide application in nonlinear mechanics [1-3]. In this paper, the results obtained are used to prove the asymptotic stability of periodic motion of the vibroimpact system described in [4].

Let $x(t)$ be a solution of (1) [or (2)], defined on the interval I (I can be the real semiaxis or line).

We will say that a solution $y(t)$ of this equation lies in an ϵ -neighborhood of a solution $x(t)$ if: 1) the measure of the symmetric difference of the domains of existence of x and y is not larger than ϵ ; 2) the points of discontinuity of y are located in an ϵ -neighborhood of the points of discontinuity of x ; 3) for all t lying outside ϵ -neighborhoods of the points of discontinuity of $x(t)$, we have $\|y(t) - x(t)\| < \epsilon$.

We will call the topology defined using ϵ -neighborhoods a B -topology.

We will call a solution $x(t)$ B -stable if it is defined for $t \geq t_0$, and for any $\epsilon > 0$ there is $\delta > 0$ such that a solution $y(t)$ satisfying $\|y(t_0) - x(t_0)\| < \delta$ belongs to an ϵ -neighborhood of $x(t)$.

A B -stable solution is B -asymptotically stable if there exists $\delta > 0$ such that for any $\epsilon > 0$ there is a number $\theta > t_0$ such that a solution y for which $\|y(t_0) - x(t_0)\| < \delta$, belongs to an ϵ -neighborhood of $x(t)$ for $t \geq \theta$.

One can verify that a solution of (2) depends continuously in the B -topology on the initial data t_0 and x_0 and the parameter μ in any compactum from \mathbb{R}^{1+n+m} , containing (t_0, x_0, μ) only if $t_0 \neq t_i(x_0, \mu)$ for all $i \in \mathbb{Z}$.

Let us now assume that in the domain

$$F = \{(x, t, \mu) \mid \|x - x_0\| < d, \quad t_0 < t < t_0 + T, \quad \|\mu - \mu_0\| < b\}$$

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the functions $f(t, x, \mu)$, $I_j(x, \mu)$, $t_j(x, \mu)$ have continuous and bounded derivatives with respect to x_j and μ_k , $j = \overline{1, n}$, $k = \overline{1, m}$

Let the solution $x(t) = x(t, \mu_0)$, $x(t_0) = x_0$ of (2) be defined for all $t \in (t_0, t_0 + T)$. The solution $y(t) = x(t, \mu_0 + \Delta\mu_1)$, $\Delta\mu_1 = (\Delta\mu, 0, \dots, 0)$ also satisfies the initial condition $y(t_0) = x_0$.

Let us call a piecewise-continuous function $u(t)$ a B-derivative of $x(t)$ with respect to μ_1 if $y(t)$ is located in a θ -neighborhood of $x(t) + \Delta\mu u(t)$, $\theta = o(\Delta\mu)$ and, furthermore, for all t , located outside θ -neighborhoods of the points of discontinuity of $x(t)$, we have $\|y(t) - x(t) - \Delta\mu u(t)\| < \nu$, $\nu = o(\Delta\mu)$.

The B-derivatives with respect to all the remaining μ_k , $k = \overline{2, m}$, are similarly defined.

Let τ_i , $i = \overline{1, p}$ be the points of discontinuity of $x(t)$ belonging to $(t_0, t_0 + T)$. Let us denote

$$A(t) = \frac{\partial f(t, x(t), \mu_0)}{\partial x}, \quad P_i = V_i - W_i + \frac{\partial I_i(x(\tau_i), \mu_0)}{\partial x} (E + V_i),$$

$$g_j(t) = \frac{\partial f(t, x(t), \mu_0)}{\partial \mu_j},$$

$$J'_i = \frac{\frac{\partial I_i(x(\tau_i), \mu_0)}{\partial \mu_j} (f(\tau_i, x(\tau_i), \mu_0) - f(\tau_i, x(\tau_i +), \mu_0))}{1 - \langle \frac{\partial I_i(x(\tau_i), \mu_0)}{\partial x}, f(\tau_i, x(\tau_i), \mu_0) \rangle} + \frac{\partial I_i(x(\tau_i), \mu_0)}{\partial \mu_j}.$$

The matrices V_i and W_i are such that for any $z \in \mathbb{R}^n$

$$V_i z = \frac{\langle \frac{\partial I_i(x(\tau_i), \mu_0)}{\partial x}, z \rangle f(\tau_i, x(\tau_i), \mu_0)}{1 - \langle \frac{\partial I_i(x(\tau_i), \mu_0)}{\partial x}, f(\tau_i, x(\tau_i), \mu_0) \rangle},$$

$$W_i z = \frac{\langle \frac{\partial I_i(x(\tau_i), \mu_0)}{\partial x}, z \rangle f(\tau_i, x(\tau_i +), \mu_0)}{1 - \langle \frac{\partial I_i(x(\tau_i), \mu_0)}{\partial x}, f(\tau_i, x(\tau_i), \mu_0) \rangle}.$$

The scalar product in $\langle \cdot, \cdot \rangle$ is denoted by the angular brackets

THEOREM 1. If system (2) satisfies all the conditions enumerated above, then for each μ_j , $j = \overline{1, m}$, there exists a B-derivative of $x(t)$ which is a solution of a linear system of differential equations with the impulse action

$$\frac{du}{dt} = A(t)u + g_j(t), \quad t \neq \tau_i, \quad \Delta u|_{t=\tau_i} = P_i u + J'_i \quad (4)$$

with initial condition $u(t_0) = 0$.

Proof. Let us verify the theorem for $j = 1$. Let θ_i be the points of discontinuity of $y(t)$. For simplicity, without destroying generality, we will let $\theta_i \geq \tau_i$ for all i .

Let us denote $v(t) = y(t) - x(t) - \Delta\mu u(t)$, where $u(t)$ is a solution of (4) for $j = 1$ with $u(t_0) = 0$.

The solutions $x(t)$, $y(t)$, $u(t)$ have the integral representations

$$x(t) = x_0 + \int_{t_0}^t f(\tau, x(\tau), \mu_0) d\tau + \sum_{t_0 < \tau_i < t} I_i(x(\tau_i), \mu_0),$$

$$y(t) = x_0 + \int_{t_0}^t f(\tau, y(\tau), \mu_0) d\tau + \sum_{t_0 < \theta_i < t} I_i(y(\theta_i), \mu),$$

$$u(t) = \int_{t_0}^t (A(\tau)u(\tau) + g_1(\tau)) d\tau + \sum_{t_0 < \tau_i < t} (P_i u(\tau_i) + J_i) \quad (5)$$

Using (5) and the differentiability of $f(t, x, \mu)$, $I_1(x, u)$, and $t_i(x, \mu)$, we find that for all $t \in (\tau_i, \theta_i)$

$$\|v(t)\| \leq e^{kT} (\rho\omega(\Delta\mu) + \Delta\mu \xi(\Delta\mu) T) \quad (6)$$

where $k = \sup_{t_0 \leq t \leq t_0+T} \|A(t)\|$, $\omega = o(\Delta\mu)$, $\xi(\Delta\mu) = o(\Delta\mu)$.

Furthermore, for all $\Delta t_i = \theta_i - \tau_i$, $i = \overline{1, p}$, we have

$$\Delta t_i = \frac{\Delta\mu \left(\left\langle \frac{\partial t_i(x(\tau_i), \mu_0)}{\partial x}, u(\tau_i) \right\rangle + \frac{\partial t_i(x(\tau_i), \mu_0)}{\partial \mu_i} \right)}{1 - \left\langle \frac{\partial t_i(x(\tau_i), \mu_0)}{\partial x}, f(\tau_i, x(\tau_i), \mu_0) \right\rangle} + o(\Delta\mu) \quad (7)$$

Relations (6) and (7) prove the theorem for $j = 1$. Its validity is verified in the same way for all remaining $j = \overline{2, m}$. The theorem is proved.

Now let us assume that in (1) the functions $f(t, x)$, $I_1(x)$, and $t_i(x)$ are continuously differentiable in x_k , $k = \overline{1, n}$.

Let Eq. (1) have a family of T -periodic solutions $x(t, \gamma)$ depending on γ in some compactum $K \subset \mathbb{R}^n$, $m < n$, inside K for each γ_j , $j = \overline{1, m}$, and B -derivatives $\varphi_j(t) = \partial x(t, \gamma) / \partial \varepsilon_j$. Let us assume that these derivatives are linearly independent for $t = 0$. Without destroying generality, we can consider that each point $x(0, \gamma)$ does not belong to one of the surfaces of discontinuity, along with some neighborhood of it.

For an arbitrary solution $x(t)$ of (1) let us define $z(t) = x(t) - x(t, \gamma)$.

Let us denote by $\tau_i = \tau_i(\gamma)$ the points of discontinuity of $x(t, \gamma)$, and by θ_i the points of discontinuity of $x(t)$. Let $\theta_i \geq \tau_i$, $i = \overline{1, p}$.

Applying the method of proving Theorem 1, one can show that for $t \in (\tau_i, \theta_i)$ the function $z(t)$ coincides with the solution of

$$du/dt = A(t)u + g(t, u), \quad t \neq \tau_i, \quad \Delta u|_{t=\tau_i} = P_i u + J_i(u), \quad (8)$$

in which the matrices $A(t)$ and P_i and the functions $g(t, u)$ and $J_i(u)$ depend on γ . The matrices $A(t)$ and P_i are determined by formulas (3), where $x(t) = x(t, \gamma)$. The conditions $\|g(t, u)\| = o(\|u\|)$ and $\|J_i(u)\| = o(\|u\|)$ are satisfied,

The next proposition is the analog of a theorem from [6].

THEOREM 2. For each $\gamma \in K$ among the characteristic exponents of the system

$$dv/dt = A(t)v, \quad t \neq \tau_i, \quad \Delta v|_{t=\tau_i} = P_i v \quad (9)$$

let there be $n - m$, having negative real parts.

Then there exists a real manifold $W(\gamma)$ of dimension $n - m$, containing a point $x(0, \gamma)$ such that a T -periodic solution $x(t, \gamma)$ is conditionally B -asymptotically stable with respect to $W(\gamma)$.

Proof. By Theorem 1, the B -derivatives φ_j are T -periodic solutions of (9). Since the functions φ_j are linearly independent for $t = 0$, they are linearly independent solutions of this system [5]. Therefore, Eqs. (9) have a characteristic exponent equal to zero, with a multiplicity not equal precisely to m .

Let us denote by $Y(t)$ the (normed at zero) fundamental matrix of the solutions of system (9), and C the matrix for which

$$Y(t+T) = Y(t)C. \quad (10)$$

Without destroying the generality of subsequent arguments, one can consider that

$$C = e^{BT}, \quad (11)$$

where $B = \text{diag}[0, B_1]$, B_1 is a matrix of dimension $(n - m) \times (n - m)$, all the eigenvalues of which have negative real parts. Let us denote $Z(t) = Y(t)e^{-Bt}$, $V_1(t, s) = Z(t) \text{diag}[0, e^{B_1(t-s)}]Z^{-1}(s)$, $V_2(t, s) = Z(t) \text{diag}[E, 0,]Z^{-1}(s)$. It follows from the expressions for V_1 and V_2 that there exists constants $M > 1$ and $\alpha < 0$, for which for $t > s$

$$\|V_1(t, s)\| \leq Me^{\alpha(t-s)}, \quad \|V_2(t, s)\| \leq M. \quad (12)$$

One can verify that the solution of the integral equation

$$\zeta(t) = Y(t)a + \int_0^t V_1(t, s)g(\zeta, s)ds - \int_t^\infty V_2(t, s)g(\zeta, s)ds + \sum_{t < \tau_i < t} V_1(t, \tau_i)J_i(\zeta(\tau_i)) - \sum_{t < \tau_i} V_2(t, \tau_i)J_i(\zeta(\tau_i)) \quad (13)$$

is also a solution of (8).

With the help of (12) and the properties of $g(t, x)$ and $J_i(x)$ it is proved by the method of sequential approximations that for sufficiently small h , $\|a\| < h$, there exists a solution $\zeta(t)$ of (13) for which

$$\|\zeta(t)\| < M_1 \exp(\alpha t/2)$$

uniformly with respect to t and a , $t > 0$, $\|a\| < h$, $0 < M_1 < +\infty$.

Substituting in (13) $t = 0$, we obtain that $\zeta_i(0) = a_i$ if $i > m$, and $\zeta_i = \Phi_i(a_{m+1}, a_{m+2}, \dots, a_n)$, $i = \overline{1, m}$, where $\Phi_i(a_{m+1}, \dots, a_n) = o(\|a\|)$. This means

$$\zeta_i = \Phi_i(\zeta_{m+1}, \zeta_{m+2}, \dots, \zeta_n) \quad (14)$$

and

$$\Phi_i(\zeta_{m+1}, \zeta_{m+2}, \dots, \zeta_n) = o(\|\zeta\|). \quad (15)$$

The relations (14) determine a manifold $W(\gamma)$, each point $(\zeta_1, \zeta_2, \dots, \zeta_n)$ of which, situated in a sufficiently small neighborhood of the origin of the coordinates, is an initial value of a solution of (8), approaching zero as $t \rightarrow +\infty$. Hence, the proof of the assertion of the theorem follows from the continuous dependence in the B-topology of the solutions of (1) on the initial data.

A corollary of Theorem 2 is an analog of the Lyapunov-Poincaré theorem on the asymptotic stability of a unique T-periodic solution of (1).

Let us consider the dynamic system analyzed in [4]. It consists of a globule bouncing on a platform moving vertically. It is assumed that the platform does not respond to the impacts of the globule and moves according to the law $X = X_0 \sin \omega t$.

The motion of the globule between kicks is determined by the formula $x = 1/2gt^2 + x^0t + x^0$. It was proved in [4] that under the condition

$$\omega^2 > \frac{\pi g}{X_0} \frac{1-R}{1+R}, \quad (16)$$

where R is the restoration coefficient, such a dynamic model admits two periodic motions $x_1(t)$ and $x_2(t)$ with a period equal to the period of oscillations of the platform.

It is assumed that the periodic motions experience an impulse action once in the period, at the moments $t = \varphi_1$ and $t = \varphi_2$, $0 < \varphi_1, \varphi_2 < 2\pi/\omega$, respectively.

Using the results of [4], we find that the mathematical model for analyzing a system on $[0, 2\pi/\omega]$ has the form of differential equations with the impulse action

$$\begin{aligned} dx_1/dt = x_2, \quad dx_2/dt = -g, \quad t \neq \frac{1}{\omega} \arcsin \frac{x_1}{X_0} \equiv t_0(x), \quad \Delta x_1|_{t=t_0(x_1)} = 0, \\ \Delta x_2|_{t=t_0(x_1)} = (1+R) \left[X_0 \omega \cos \left(\arcsin \frac{x_1}{X_0} - x^0 + \frac{g}{\omega} \arcsin \frac{x_1}{X_0} \right) \right], \end{aligned} \quad (17)$$

where \dot{x}^0 is the velocity at $t = 0$.

The systems of linear equations

$$du_1/dt = u_2, \quad du_2/dt = -g, \quad t \neq \varphi_1, \quad \Delta u_1|_{t=\varphi_1} = 0, \quad \Delta u_2|_{t=\varphi_1} = b_+ u_1, \quad (18)$$

and

$$du_1/dt = u_2, \quad du_2/dt = -g, \quad t \neq \varphi_2, \quad \Delta u_1|_{t=\varphi_2} = 0, \quad \Delta u_2|_{t=\varphi_2} = b_- u_1, \quad (19)$$

where

$$b_+ = \frac{(1+R)^2 \omega}{\pi(1-R)} - (1+R)\omega \sqrt{\left[\frac{X_0 \omega^2 (1+R)}{\pi g (1-R)} \right]^2 - 1},$$

$$b_- = \frac{(1+R)^2 \omega}{\pi(1-R)} + (1+R)\omega \sqrt{\left[\frac{X_0 \omega^2 (1+R)}{\pi g (1-R)} \right]^2 - 1}$$

are, respectively, systems of equations in variations for equations (17) with respect to $x_1(t)$ and $x_2(t)$.

The characteristic equation [5] for Eqs. (18) and (19), respectively, has the form

$$\lambda^2 - 2 \left(1 + \frac{\pi b_+}{\omega} \right) \lambda + 1 = 0 \quad (20)$$

and

$$\lambda^2 - 2 \left(1 + \frac{\pi b_-}{\omega} \right) \lambda + 1 = 0. \quad (21)$$

Solving these equations, we obtain that Eq. (21) does not have roots with a nonpositive real part. Consequently, the solution $x_2(t)$ is unstable.

Both roots of (20) have negative real parts if

$$\omega^2 > \sqrt{\left[\frac{\pi g (1-R)}{X_0 (1+R)} \right]^2 + \left[\frac{g}{X_0} \left(1 + \frac{1-R}{(1+R)^2} \right) \right]^2}. \quad (22)$$

Comparing (16) and (22), we find that by the analog of the Lyapunov-Poincaré theorem on the stability of a periodic solution, if (22) is valid, the vibroimpact system under investigation admits a unique B-asymptotically stable $2\pi/\omega$ -periodic solution.

The result obtained agrees well with the deduction of [4].

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